# Interacting Electron-Photon System 

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## Outline

(1) Introduction
(2) Single Photon System
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4 Two-body, Non-interacting System
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## Introduction

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- We conducted our research gradually, first examining the system of a single photon, then that of a single electron, then that of the two without any interaction, and, at last, that of two interacting.


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- In one space dimension this means:

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- Kiessling \& Tahvildar-Zadeh also discovered this equation in 2018. It is a Dirac-type equation, and in 1-dim. it reads:

$$
-i \hbar \gamma^{\mu} \frac{\partial \Psi_{p h}}{\partial x^{\mu}}=0
$$

where $\hbar=$ reduced Planck's constant, $x^{0}=t, x^{1}=s$,
$\gamma^{0}=\left(\begin{array}{ll}0 & 1 \\ 1 & 0\end{array}\right), \gamma^{1}=\left(\begin{array}{ll}0 & -1 \\ 1 & 0\end{array}\right)$, and repeated indices are summed over the range $\mu=0,1$.

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## Photon Probability Current and Velocity Field

- According to Kiessling and Tahvildar-Zadeh, in one space dimension the quantum probability current of detecting the photon is

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j_{p h}^{\mu}(\text { time }, \text { position })=\frac{1}{4} \operatorname{trace}\left(\overline{\Psi_{p h}} \gamma^{\mu} \Psi_{p h} \gamma(X)\right)
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- Here $X=\left(X^{0}, X^{1}\right)$ is a constant vector field computed from $\Psi_{p h}^{0}$, $\gamma(X):=\gamma_{0} X^{0}+\gamma_{1} X^{1}$, and $\bar{\Psi}:=\gamma^{0} \psi^{\dagger} \gamma^{0}$.


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- $\rho(t, s)$ is normalized, namely $\int \rho(t, s) d s=1$.


## Photon Probability Density

- The photon probability density looks like this: http://reu.dimacs.rutgers.edu/~aas377/photon_pdf.mp4


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- The photon probability density looks like this: http://reu.dimacs.rutgers.edu/~aas377/photon_pdf.mp4
- Varying initial conditions gives us the following: http://reu. dimacs.rutgers.edu/~aas377/multiple_photon_pdf.mp4


## The Guiding Equation

- The motion of the photon is guided by its wave function:

$$
\left\{\begin{array}{l}
\frac{d q}{d t}=v_{p h}(t, q(t))=\frac{j^{1}(t, q(t))}{j^{0}(t, q(t))} \\
q(0)=q_{0}
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- $q_{0}$ is the actual initial position of the photon. All we know about it is that it is randomly distributed according to the initial probability density $\rho(0, s)$.



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- Like in the case of a single photon, the wave function of a single electron also satisfies a relativistic equation. In particular, it satisfies the massive Dirac equation:

$$
-i \hbar \gamma^{\mu} \partial_{\mu} \Psi_{e l}+m_{e l} \Psi_{e l}=0
$$

where $m_{e l}=$ the mass of electron.

## Electron Probability Current and Velocity Field, and Guiding Equation

- The probability current of an electron is known:

$$
j_{e l}^{\mu}(\text { time }, \text { position })=\overline{\Psi_{e l}} \gamma^{\mu} \Psi_{e l}
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- Similarly to the photon case, the guiding equation for the electron is:

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\frac{d q}{d t}=v_{e l}(t, q(t))=\frac{j^{1}(t, q(t))}{j^{0}(t, q(t))} \\
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## Electron Probability Density

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## Electron Trajectories and Parameters

- The electron wave function $\psi_{e l}$, has a mass term: $\omega=$ mass $/ \hbar$, and a parameter we can change: standard deviation of the initial distribution: $\sigma$. The following graph shows the trajectory of an electron guided by the velocity field with different parameters.


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## Two-body, Non-interacting System

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- The wave function, $\psi$, is a function of four variables, namely the time and position of each particle.
- To get a wave function that describes both a photon and an electron in a non-interacting system, we take the Tensor Product $(\otimes)$ of the electron and the photon wave functions, giving us a four component object $\psi=\left(\psi_{++}, \psi_{+-}, \psi_{-+}, \psi_{--}\right)$


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- The guiding equations for photon and electron are derived using the Hypersurface Bohm-Dirac (HBD) Theory, which allows us to describe the motion of the photon and electron in a common time


## Wave Equation and Probability Current

- The tensored wave function satisfies a relativistic wave equation obtained by Tensor Product of the photon and electron wave equations

$$
\left\{\begin{array}{l}
-i \hbar \gamma^{\mu} \partial x_{p h}^{\mu} \psi=0 \\
-i \hbar \gamma^{\mu} \partial x_{e l}^{\mu} \psi+m_{e l} \psi=0 \\
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- Since Tensor Product preserves probability distributions, "multiplying" the probability density movies of the photon and electron gives the joint probability density for the non-interacting system: http: //reu.dimacs.rutgers.edu/~aas377/non_interacting2.mp4


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- The probability current is the following:

$$
\left(\begin{array}{cc}
j^{10}=\left|\psi_{++}\right|^{2}+\left|\psi_{+-}\right|^{2}-\left|\psi_{-+}\right|^{2}-\left|\psi_{--}\right|^{2} \\
j^{01} & =\left|\psi_{++}\right|^{2}-\left|\psi_{+-}\right|^{2}+\left|\psi_{-+}\right|^{2}-\left|\psi_{--}\right|^{2} \\
j^{00} & =\left|\psi_{++}\right|^{2}+\left|\psi_{+-}\right|^{2}+\left|\psi_{-+}\right|^{2}+\left|\psi_{--}\right|^{2} \\
j^{11} & =\left|\psi_{++}\right|^{2}-\left|\psi_{+-}\right|^{2}-\left|\psi_{-+}\right|^{2}+\left|\psi_{--}\right|^{2}
\end{array}\right)
$$

## Trajectories of the Two-body, Non-interacting System

- The following graph shows the trajectories of a non-interacting system of one electron and one photon.



## Chapter 4: Two-body, Interacting System

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## Chapter 4: Two-body, Interacting System

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- We do this by adding a boundary condition: we set the relative velocities of photon and electron to be 0 when the particles are at the same space and time.
- Adding the boundary condition to the wave function gives us a modified probability density function: http:
//reu.dimacs.rutgers.edu/~aas377/interacting_pdf_2.mp4


## Trajectories of the Two-body, Interacting System

- Adding the boundary condition to the wave function gives us the trajectories of an interacting electron-photon system.



## Varying Parameters in the Interacting System

- Changing the sigma and omega of the electron gives us the following changes in trajectories of the electron:



## Varying Parameters in the Interacting System

- Changing the polarization angles of electron and photon gives us:











## Varying Parameters in the Interacting System

- Changing the mean momentum of the incoming photon varies the trajectories as follows:



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- Changing the mean momentum of the incoming photon varies the trajectories as follows:



- Momentum is related to energy, so if the photon does not have enough momentum, it cannot get the electron to bounce away.


## Thank You

- This research is made possible by the Rutgers Math Department, REU-DIMACS, and the generous help from Professor Shadi Tahvildar-Zadeh

